Electronic and magnetic structure of the \((\text{LaMnO}_3)_{2n}/(\text{SrMnO}_3)_n\) superlattices

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We study the magnetic structure of the \((\text{LaMnO}_3)_{2n}/(\text{SrMnO}_3)_n\) superlattices from density-functional calculations. In agreement with the experiments, we find that the magnetism changes with the layer thickness \(n\). The reason for the different magnetic structures is shown to be the varying potential barrier across the interface, which controls the leakage of the Mn-\(e_g\) electrons from the LaMnO\(_3\) side to the SrMnO\(_3\) side. This in turn affects the interfacial magnetism via the carrier-mediated Zener double exchange. For the \(n=1\) superlattice, the Mn-\(e_g\) electrons are more or less spread over the entire lattice so that the magnetic behavior is similar to the equivalent alloy compound La\(_{2}\)Sr\(_{1/3}\)MnO\(_3\). For larger \(n\), the \(e_g\) electron transfer occurs mostly between the two layers adjacent to the interface, thus leaving the magnetism unchanged and bulklike away from the interface region.

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I. INTRODUCTION

Superlattices made up of strongly correlated transition-metal oxides such as LaMnO\(_3\) (LMO) and SrMnO\(_3\) (SMO) are of current interest because of the diverse magnetic and electronic phases they exhibit. For example recent experimental results reveal that (LMO)\(_{2n}/(\text{SMO})_n\) superlattice is uniformly ferromagnetic for the short-period structure \((n=1)\), while the long-period superlattices \((n \geq 3)\) show bulk antiferromagnetic ordering away from the interface and ferromagnetic ordering at the interface.\(^{1,2}\)

In this paper, we report results of our electronic structure calculations, based on the density-functional theory (DFT), performed to understand the change in the magnetic properties of the (LMO)\(_{2n}/(\text{SMO})_n\) superlattices as a function of the layer thickness \(n\). We show that there exists a potential barrier for the electrons, in particular, for the Mn-\(e_g\) electrons, the strength of which differs with the layer thickness \(n\). This varying potential barrier, which controls the leakage of the Mn-\(e_g\) electrons from the LMO side to SMO side, in turn determines the stable magnetic configurations in the (LMO)\(_{2n}/(\text{SMO})_n\) superlattices. In agreement with the experiments, our calculations predict a uniform ferromagnetic (FM) ordering in the short-period superlattice \((n=1)\) and the co-existence of interface FM phase and inner bulk antiferromagnetic (AFM) phases in the long-period superlattices \((n \geq 3)\). The magnetism can be qualitatively understood in terms of the two competing interactions, viz., the antiferromagnetic superexchange between the core spins and the Zener ferromagnetic double exchange mediated by the itinerant \(e_g\) electrons.

II. COMPUTATIONAL AND STRUCTURAL DETAILS

The results presented in this paper are obtained from the DFT studies of three superlattices, namely, (LMO)\(_2/(\text{SMO})_1\) (schematically shown in Fig. 1), (LMO)\(_3/(\text{SMO})_2\), and (LMO)\(_6/(\text{SMO})_3\) using the linear muffin-tin orbitals (LMTO) method\(^3\) with general gradient approximation\(^4\) and on-site Coulomb correction (GGA+U)\(^5\). The Coulomb (\(U\)) and the exchange parameter (\(J\)) are taken as 5 and 1 eV, respectively. Each superlattice consists of twice the formula unit because of the magnetic structures considered in the paper. The bulk lattice parameters of LMO and SMO are, respectively, 3.935 and 3.802 Å. However, since most of the experimental results reported in the literature are based on the LMO/SMO superlattices grown on the SrTiO\(_3\) (STO) substrate,\(^1,2,6,7\) we have taken the in-plane lattice parameter for the (LMO)\(_{2n}/(\text{SMO})_n\) superlattices as the bulk STO lattice parameter (3.905 Å). The out-of-plane lattice parameters are taken to be 3.99 (LMO) and 3.65 Å (SMO) which preserve the bulk volumes. A somewhat better estimate of the out-of-plane lattice parameters may be obtained from the linear relation containing the Poisson’s ratio,\(^8–10\) which would yield the values 3.95 (LMO) and 3.78 Å (SMO). We do not expect these differences to change the basic physics discussed here. However, a substantial change in the strain condition, obtained for example by growing the superlattice on different substrates, can alter the orbital ordering and through it the interfacial magnetic structure as discussed below.

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![FIG. 1. Schematic unit cell of (LMO)\(_2/(\text{SMO})_1\) superlattice and the magnetic structure as predicted from the DFT calculations. Mn-0 represents the interfacial Mn atoms surrounded by both SrO and LaO layers and Mn-1 represents the Mn atoms inside the LMO part. Because the SMO part is small, there is no Mn atom surrounded by two SrO layers in this structure. The nearest-neighbor Mn-Mn exchange interactions are indicated by the J’s.](image-url)
elsewhere. The basal Jahn-Teller (JT) distortion \((Q_2)\) for the inner Mn layers in the LMO site is taken the same as the bulk value (0.15 Å). The value of \(Q_2\) for the interface Mn layers is taken as 0.07 Å in view of the fact that the JT distortion is reduced in the mixed compound \((\text{La}, \text{Sr})\text{MnO}_3\), and one expects the distortion to scale roughly linearly with the number of \(e_g\) electrons on the Mn atom, which is approximately half for the interfacing Mn atom.

### III. ELECTRONIC STRUCTURE OF THE \((\text{LMO})_2/(\text{SMO})_1\) SUPERLATTICE

Before discussing the electronic and magnetic properties of the \((\text{LMO})_2/(\text{SMO})_n\) superlattices, we summarize the electronic structure and magnetism of the bulk SMO and LMO compounds. In bulk SMO, the Mn atoms are in the 4+ charged state so that they have three \(d\) electrons occupying the triply-degenerate \(t_{2g}\) states. The doubly-degenerate \(e_g\) states, which are higher in energy with respect to the \(t_{2g}\) states because of the MnO\(_6\) octahedral crystal field, remain unoccupied. The \(t_{2g}\) core spins interact via an antiferromagnetic superexchange so as to stabilize the \(G\)-type AFM ordering in the bulk SMO compound.\(^{11,12}\)

In bulk LMO, the Mn atoms are in the 3+ charged state with four occupied \(d\) electrons. Three electrons are present in the \(t_{2g}\) states and the remaining one in the \(e_g\) states. The Jahn-Teller distortion of the MnO\(_6\) octahedron further splits the \(e_g\) states into two nondegenerate states: \(e_g^1\), which is lower in energy, and \(e_g^2\), which is higher in energy.\(^{13}\) The \(e_g^1\) orbital, occupied by the lone electron, has its lobes pointed toward the longest Mn-O bond. The JT distortion stabilizes the \(A\)-type AFM structure in the LMO compound due to a combination of the superexchange and Zener double exchange.\(^{14}\)

The charge reconstruction at the LMO/SMO interface\(^{15,16}\) is expected to change the electronic and magnetic properties of the \((\text{LMO})_2/(\text{SMO})_n\) superlattices, which will be discussed in the remaining part of the paper.

Out of a number of magnetic configurations that we considered, the DFT calculations predict a ferromagnetic ground state for the \((\text{LMO})_2/(\text{SMO})_1\) superlattice. In Fig. 2, we have shown the total and partial spin-resolved densities of states (DOSs) for the ferromagnetic configuration of this superlattice. The characteristic features of the electronic structure as seen from the figure are as follows. The Mn-\(t_{2g}\) states lie far below the Fermi level \(E_F\) because of the octahedral crystal field and strong Coulomb repulsion. The \(O-p\) states occur in the energy range of \(-6\) to \(-1\) eV. The Mn-\(e_g\) states occur around the Fermi level \(E_F\), while the Sr-\(d\), La-\(d\), and La-\(f\) states lie far above it.

As Fig. 2 shows, the most important feature in the electronic structure of \((\text{LMO})_2/(\text{SMO})_1\) is that the delocalized \(e_g\) states of both Mn-0 (Mn atoms at the interface) and Mn-1 (Mn atoms inside the LMO part) are partially occupied, which is in agreement with the earlier electronic structure calculations.\(^{17}\) These partially occupied \(e_g\) states will mediate a strong Zener ferromagnetic double exchange between the Mn-\(t_{2g}\) core spins, which wins over the antiferromagnetic superexchange, so that a uniform ferromagnetic ordering throughout the superlattice is stabilized. The calculation of the Mn-Mn exchange interactions discussed below indicates that the FM ordering is stable, quite similar to the equivalent alloy compound La\(_{2.5}\)Sr\(_{1.5}\)MnO\(_3\).\(^{1}\)

![Figure 2](https://example.com/figure2.png)

**FIG. 2.** Total (upper panel) and partial spin-resolved DOS for the ferromagnetic \((\text{LMO})_2/(\text{SMO})_1\) superlattice. The labeling of the Mn atoms is as in Fig. 1. Upper and lower segments within each panel correspond, respectively, to the majority (↑) and minority (↓) spin densities.

### IV. MAGNETIC EXCHANGE INTERACTION

In order to study the magnetic ground state for the \((\text{LMO})_2/(\text{SMO})_n\) superlattices, we have calculated the neighboring Mn-Mn exchange interaction energies \((J's)\) for various exchange interactions as shown in Figs. 1 and 3. In these figures the symbol \(J_L\) represents the out-of-plane exchange interactions across the SrO layer close to the interface, while \(J_3\) and \(J_4\) represent the same across the LaO layers close to the interface and away from the interface, respectively. \(J_2\) denotes the in-plane exchange interaction for the interfacial MnO\(_2\) layer, which is surrounded by LaO and SrO layers, whereas \(J_3\) and \(J_4\) denote the same for the MnO\(_2\) layer inside the LMO and SMO part of the superlattice, respectively.

The exchange interaction \(J\) is defined as the energy difference between the ferromagnetic alignment and the antiferromagnetic alignment of two neighboring Mn spins \((J=E_{FM}-E_{AFM})\).
the magnetic structure as obtained from the DFT calculations. Oxygen atoms occur at the intersections of the checkerboard layers forming the MnO2 octahedron. Mn atoms of each MnO2 layer are labeled as shown in the figure. Definitions of the exchange interactions for the (LMO)4/(SMO)4 superlattice are identical to the ones shown here, and they are also consistent with Fig. 1 for the (LMO)4/(SMO)4 superlattice.

We computed them by performing a number of total-energy calculations for various magnetic configurations for each superlattice and fitting the energies with the results of a nearest-neighbor Heisenberg model. The results are listed in Table I. The in-plane magnetic interaction \( J_0 \) inside the LMO part, which is strongly ferromagnetic, was not computed. For the case of the (LMO)4/(SMO)4 superlattice, the values of the exchange interaction for the Mn layers away from the interface in the LMO part and SMO part are, respectively, 12 \( (J_4) \) and 19 meV \( (J_5) \). These values are in good agreement with the experimental results for the bulk LMO (\( J \approx 9.7 \) meV) and bulk SMO (\( J \approx 13.1 \) meV).21-23

From Table I, we see that for (LMO)2/(SMO)1, the in-plane exchange interaction \( J_1 \), as well as the out-of-plane exchange interactions \( J_2 \) and \( J_3 \), is strong and negative so as to stabilize the FM ordering throughout the superlattice, consistent with the experimental observations.1,2 Turning now to the (LMO)4/(SMO)4 superlattice, the in-plane interactions \( (J_0 \text{ and } J_2) \) are FM as also are the out-of-plane interactions \( (J_3 \text{ and } J_4) \) within the LMO part. In the SMO part, the out-of-plane exchange interaction \( (J_1) \) is AFM, but this being weaker as compared to the in-plane \( J_2 \) (FM) and \( J_3 \) (AFM), the magnetic configuration within the SMO part is controlled by the latter two exchange interactions as shown in Fig. 3. Finally, for the (LMO)6/(SMO)3 superlattice, the values of \( J \)'s are similar to those of the (LMO)4/(SMO)4 superlattice, except that now the out-of-plane exchange interaction for the inner MnO2 layers in the LMO side \( (J_4) \) is positive so as to establish an A-type AFM configuration as in the bulk LMO compound.

We note from the above discussions that as we increase the layer thickness \( n \), the FM interactions between the Mn spins occurring on the two sides of the La0 layers (see \( J_3 \) and \( J_4 \) in Table I) gradually become weak, which eventually makes the LMO part type-A AFM like in the bulk. This already happens for \( n=3 \). The transition from the FM to AFM ordering for the Mn layers away from the interface with the increase in the layer thickness \( n \) is indicative of the fact that the charge reconstruction is essentially confined to the few interface layers for the long-period superlattices \( (n \geq 3) \).

The calculated magnetic moments of the Mn atoms (within the muffin-tin sphere radius of 1.15 Å) are also given in Table I. Our magnetic unit cells consisted of two Mn atoms per layer, with slightly different magnetic moments, which we have averaged over to obtain the results presented in Table I. For the \( n=2 \) and \( 3 \) superlattices, where bulklike SMO and LMO regions exist, the magnetic moments are consistent with the \( 2^3 \uparrow \) occupancy of the Mn-1 (bulk SMO like), \( 2^3 \downarrow e_{g}^1 \uparrow \) occupancies of the Mn-1, Mn-2, Mn-3 (bulk LMO like), and the \( 2^3 \downarrow e_{g}^{0.5} \uparrow \) occupancy of the Mn-0 (interfacial Mn), which gives rise to the nominal magnetic moments of 3.5\( \mu_B \), 4\( \mu_B \), and 3.5\( \mu_B \), respectively. For the \( n=1 \) superlattice, the \( e_g \) electrons are spread more or less all over the lattice, and this is reflected in the near equality of the magnetic moments of the two Mn atoms: 3.57\( \mu_B \) for Mn-0 and 3.70\( \mu_B \) for Mn-1, as seen from Table I.

V. ELECTRIC POTENTIAL PROFILE AND CHARGE RECONSTRUCTION AT THE INTERFACE

The potential seen by the electrons varies as one crosses the interface from one side to the other. This for example

<table>
<thead>
<tr>
<th>Superlattice</th>
<th>Magnetic moment (( \mu_B ))</th>
<th>Exchange interaction (meV)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Mn-1</td>
<td>Mn-0</td>
</tr>
<tr>
<td>(LMO)4/(SMO)4</td>
<td>3.57</td>
<td>3.70</td>
</tr>
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<td>(LM)4/(SMO)2</td>
<td>2.99</td>
<td>3.52</td>
</tr>
<tr>
<td>(LMO)6/(SMO)3</td>
<td>2.98</td>
<td>3.51</td>
</tr>
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</table>
leads to the well-known band offset in the semiconductors. Our calculations show that for the present superlattices, there is a potential barrier as one goes from the LMO to the SMO side. This controls the leakage of the Mn-eg electrons across the barrier, which in turn affects the magnetic exchange interactions near the interface leading to diverse magnetic phases.

In Figs. 4(a)–4(c), we plot the calculated oxygen 1s core energies, indicating the potential barrier across the interface. However, the valence states experience a somewhat different potential than the core states because of different energy terms. Since the Mn-eg electrons are mainly the electrons that are transferred across the interface, we now examine the potential V(z) felt by these electrons. In order to obtain the variation in this potential, we have studied the band structure and the atomic characters of the wave functions in each superlattice by examining the so-called “fat” bands in the LMTO results, which indicate the relative contributions of the various orbitals to the wave function making the band. From the fat bands, the lowest Mn-eg state belonging to a particular Mn layer can be identified, which is then indicative of the potential experienced by the Mn-eg electrons in the various layers.

These results are shown in Fig. 4(d). The variation in V(z) for the n=1 superlattice is quite similar to the variation in the oxygen 1s core energies and hence is not shown in the figure. For this superlattice, we have a weakly varying potential due to the close proximity of the interfaces to one another, which results in the overlap of the attractive Coulomb potential formed by the positively charged interfacial LaO+ layers. In this case the Mn-eg electrons are more or less spread throughout the superlattice as seen from the layer-projected DOS (Fig. 2), where all Mn atoms have partially filled eg states. These itinerant eg electrons mediate the Zener double exchange stabilizing the FM ordering throughout the superlattice.

With the increase in the layer thickness n, the variation in the potential becomes stronger, leading to the formation of a potential barrier at the interface with the LMO side having a lower potential than the SMO side. This results in restricting the leakage of the Mn-eg electrons to the SMO side (Fig. 4). Thus, for example in the case of the n=3 superlattice, there is very little eg electron on the Mn-1 atom belonging to the SMO side (Fig. 5, topmost panel). Since the Mn-eg states are unoccupied in the SMO side, a G-type AFM structure is stabilized as in the bulk SMO.

The case of the n=2 superlattice is intermediate between the short-period and the long-period (n=3) superlattices. Here, on one hand, the leakage of electrons to the SMO side is small enough that the G-type AFM is maintained as there in the bulk. On the other hand the number of eg electrons leaving the LMO side is large enough that the LMO part behaves like a hole doped bulk (La1-xSrxMnO3), thereby stabilizing the FM structure as in the short-period superlattice (n=1). However, as the calculated ferromagnetic stabilization energy is relatively small here as compared to the n=1 case, it is only weakly ferromagnetic (JF =−4 meV, Table I).

In contrast to this, in the long-period superlattices (n ≥3), a much stronger potential barrier prevents any significant leakage of the electrons to the SMO side, except to the very first interfacial layer. This leads to the bulk magnetic behavior inside the LMO as well as the SMO parts. The only layers affected by the electron leakage are just two layers at the interface so that the magnetic structure as indicated in Fig. 4(a) is of the type ...[FGGF-FAAAF], where the vertical line indicates the interface. The calculated ground-state magnetic structures for the three superlattices discussed in this paper agree with those observed in the experiments.1,2

VI. ELECTRONIC STRUCTURE OF THE (LMO)6/(SMO)3 SUPERLATTICE

We now turn to the electronic structure of the (LMO)6/(SMO)3 superlattice, which would be typical of the
As seen from the figure, deep inside both the SMO and LMO parts, the electron occupancies are more or less similar to those of the respective bulk compounds. The bulk behavior occurs already beyond just one Mn layer on either side of the interface. In the SMO part (Mn-1 densities, topmost panel), the Mn-\(t_{2g}\) spin-up states are filled while the \(e_g\) states are empty just like bulk SMO. In the LMO part (Mn-2 densities, bottommost panel), the \(e_g\) states are Jahn-Teller split into two bands, with the lower one occupied, again, as in the bulk LMO. As one approaches the interface from the LMO side, the \(e_g\) occupancy is reduced slightly from one due to the leakage of the electron to the interfacial Mn-0 layer. The transferred electron across the interface controls the magnetic behavior of the interfacial layers as already discussed.

**SUMMARY**

In summary, we have studied the change in the magnetic properties of the (LMO)\(_n\)/SMO\(_n\) superlattices as a function of the layer thickness \(n\) and explained the observed magnetic structure in terms of the electron leakage across the interface and a double exchange interaction between these electrons and the Mn-\(t_{2g}\) moments. For the short-period superlattice (\(n=1\)), we find a weak variation in the potential leading to the spreading of the Mn-\(e_g\) electrons throughout the superlattice, resulting in a FM structure via the carrier-mediated Zener double exchange, much like the alloy compound La\(_{3}\)Sr\(_{1}\)MnO\(_3\). For higher \(n\) there is a potential barrier restricting the electron leakage to the SMO side. For \(n \geq 3\), the charge leakage is restricted to just two layers at the interface, beyond which a bulklike electronic and magnetic structure results.

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FIG. 5. Spin-up (↑) and spin-down (↓) Mn-\(d\) DOS for the \(n=3\) superlattice. Up and down spins are with respect to the local magnetic moment of the Mn atom. The labeling of the Mn atoms is the same as in Fig. 4. The projected Mn-3 densities (not shown here) are similar to the Mn-2 densities as the bulk limit has already been reached.


C. Zener, Phys. Rev. 82, 403 (1951).


